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FURTHER GENERALIZATION OF KOBAYASHI'S GAMMA *k*-FUNCTIONS

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ABSTRACT. In this present paper, we deal with the more generalized functions involving Gauss hypergeometric k-functions $F_k((a; k), (b; k); (c, k); z)$

$$\begin{split} D_k \left(\begin{array}{c} (a;k), (b;k); (c,k), p \\ u, v, \delta \end{array} \right) \\ &= v^{\frac{-a}{k}} \int_0^\infty t^{u-1} (1-k\frac{t}{v})^{\frac{\delta}{k}-1} e^{-\frac{pt^k}{k}} F_k((a;k), (b;k); (c,k); -\frac{t}{v}) dt, \end{split}$$

where $|\arg v| < \pi$, Re(u) > 0. This function reduces to a generalized gamma kfunction recently defined by "[Mubeen et al.; Jour. Ineq. Spec. Funct. Volume 7 Issue 1(2016), pp.53-60]" when $\delta = k$. Also, it reduces to gamma k-function Γ_k when p = 1, $\delta = k$ and b = c. In this paper, we obtain several properties for the newly defined generalized function D_k , its recurrence relations and asymptotic expansion for D_k where k > 0. We establish further generalization of Kobayashi's gamma k-function i.e., $\Gamma_{m,k}(u, v)$, which is an improved version of Kobayashi's gamma function.

1. INTRODUCTION

Kobayashi [11] has considered plane wave diffraction theory by a strip using the Wiener-Hopf techniques and further he has established a complete high frequency asymptotic solution which is uniformly valid everywhere in space and which has no limitations on incidence and observation angles. The asymptotic solutions established by him can predict a high frequency field behavior everywhere in space. Whereas dealing with the asymptotic expansion of a class of branch-cut integrals arises in diffraction theory related to Wiener-Hopf techniques, Kobayashi [12] also established integrals of the following form

$$\Gamma_m(u,v) = \int_0^\infty \frac{t^{u-1}e^{-t}}{(t+v)^m} dt, \quad |\arg v| < \pi, Re(u) > 0, \tag{1.1}$$

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where $m \in \mathbb{Z}^+$, the function $\Gamma_m(u, v)$ is the generalized gamma function. When m = 0, (1.1) reduces to the well known gamma function

$$\Gamma(u) = \int_{0}^{\infty} t^{u-1} e^{-t} dt, Re(u) > 0$$
(1.2)

and he proved some analytical properties of $\Gamma_m(u, v)$ including regularity, asymptotic expansion and analytic continuation.

In (1998), Al-Musallam and Kalla [3] defined a more general function involving Gauss hypergeometric function F(a, b; c; z) which is given by

$$D\left(\begin{array}{c}a,b;c,p\\u,v\end{array}\right) = v^{-a} \int_{0}^{\infty} t^{u-1} e^{-pt} F(a,b;c,-\frac{t}{v}) dt, |\arg v| < \pi, Re(u) > 0.$$
(1.3)

In the same paper, they investigated several properties for the function D including analytical continuation, recurrence formulae and computation for special values of parameters and some new properties of the generalized gamma function $\Gamma_m(u, v)$. For further detail about the function D, see ([1, 2, 5],[7],[8],[10],[13, 14, 17]).

In (2001), Galue *et al.* [9], have introduced further generalization of Kobayashi's gamma function and they defined

$$D\left(\begin{array}{c}a,b;c,p\\u,v,\delta\end{array}\right) = v^{-a} \int_{0}^{\infty} t^{u-1} (1-\frac{t}{v})^{\delta-1} e^{-pt} F(a,b;c,-\frac{t}{v}) dt,$$
(1.4)

where $|\arg v| < \pi$, Re(p) > 0 and Re(u) > 0. This reduces to Kobayashi's [11] generalized gamma function when $\delta = 1$, p = 1 and b = c. Also, it reduces to a more generalized function defined by the researchers [3, 4]. In the same paper, they have defined the generalized incomplete and the complementary incomplete functions associated with $D\begin{pmatrix} a, b; c, p \\ u, v, \delta \end{pmatrix}$ and introduced their asymptotic expansions.

Diaz and Pariguan [6] introduced the generalized k-gamma function as

$$\Gamma_k(x) = \lim_{n \to \infty} \frac{n! k^n (nk)^{\frac{x}{k} - 1}}{(x)_{n,k}}, \ k > 0, x \in \mathbb{C} \setminus kZ^-$$
(1.5)

and also gave the properties of said function. The Γ_k is one parameter deformation of the classical gamma function such that $\Gamma_k \to \Gamma$ as $k \to 1$. The Γ_k is based on the repeated appearance of the expression of the following form

$$\alpha(\alpha + k)(\alpha + 2k)(\alpha + 3k)...(\alpha + (n-1)k).$$
(1.6)

The function of the variable α given by the statement (1.6), denoted by $(\alpha)_{n,k}$, is called the Pochhammer k-symbol. We obtain the usual Pochhammer symbol $(\alpha)_n$ by taking k = 1. The definition given in relation (1.5), is the generalization of $\Gamma(x)$ and the integral form of Γ_k is given by

$$\Gamma_k(x) = \int_0^\infty t^{x-1} e^{-\frac{t^k}{k}} dt, \ Re(x) > 0.$$
(1.7)

From relation (1.7), we can easily show that

$$\Gamma_k(x) = k^{\frac{x}{k} - 1} \Gamma(\frac{x}{k}). \tag{1.8}$$

Recently Mubeen *et al.* [15] have defined the following contiguous function relations for hypergeometric k-functions as

$$(a-b)F_k(a,b;c;z) = aF_k(a+k,b;c;z) - bF_k(a,b+k;c;z),$$
(1.9)

$$(a - c + k)F_k(a, b; c; z) = aF_k(a + k, b; c; z) - (c - k)F_k(a, b; c - k; z).$$
(1.10)

$$(a+b-c)F_k(a,b;c;z) = a(1-kz)F_k(a+k,b;c;z) - (c-b)F_k(a,b-k;c;z).$$
(1.11)

$$(c-a-b)F_k(a,b;c;z) = (c-a)F_k(a-k,b;c;z) - b(1-kz)F_k(a,b+k;c;z).$$
(1.12)

$$(b-a)(1-kz)F_k(a,b;c;z) = (c-a)F_k(a-k,b;c;z) - (c-b)F_k(a,b-k;c;z).$$
(1.13)

$$(b-c+k)F_k(a,b;c;z) = bF_k(a,b+k;c;z) - (c-k)F_k(a,b;c-k;z).$$
(1.14)

Very recently Mubeen *et al.* [16] have defined the generalized gamma k-function

$$D_k \left(\begin{array}{c} (a;k), (b;k); (c;k), p \\ u, v \end{array} \right) = v^{\frac{-a}{k}} \int_0^\infty t^{u-1} e^{-p\frac{t^k}{k}} F_k[(a;k), (b;k); (c;k), -\frac{t}{v}] dt,$$
(1.15)

and investigate some of its properties. In the same paper, they have defined an improved version of Kobayashi's gamma k-function of the form

$$\Gamma_{a,k}(u,v) = \int_{0}^{\infty} \frac{t^{u-1}e^{-\frac{t^{k}}{k}}}{(v+kt)^{\frac{n}{k}}} dt, \quad |\arg v| < \pi, Re(u) > 0, k > 0.$$

2. Further Generalization of Gamma k-Functions

In this section, we introduce further generalization of gamma k-functions and some of their properties where k > 0.

Definition 2.1. Let a, b, c and p be complex parameters with $c \neq 0, -1, -2, \cdots$, Re(p) > 0 and ${}_{2}F_{1,k}(a, b; c; z)$ is Gauss hypergeometric k-function, then we define

$$D_{k} \begin{pmatrix} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{pmatrix}$$

$$= v^{\frac{-a}{k}} \int_{0}^{\infty} t^{u-1} (1 - k \frac{t}{v})^{\frac{\delta}{k} - 1} e^{-p \frac{t^{k}}{k}} F_{k}[(a;k), (b;k); (c;k), -\frac{t}{v}] dt, \quad (2.1)$$

where $v \in W_v = \{v \in \mathbb{C} : |\arg v| < \pi\}, u \in W_u = \{u \in \mathbb{C} : Re(u) > 0\}$ and $F_k[(a;k), (b;k); (c;k), -\frac{t}{v}]$ is the Gauss hypergeometric k-functions. Obviously if $\delta = k$, then it reduces to $D_k \begin{pmatrix} (a;k), (b;k); (c;k), p \\ u, v \end{pmatrix}$ (see [16]). Now we prove the

generalize Kobayashi's gamma k-function $\Gamma_{m,k}(u, v)$. **Theorem 2.1.** If k > 0, then

$$\Gamma_{a,k}(u,v) = \int_{0}^{\infty} \frac{t^{u-1}e^{-\frac{t^{k}}{k}}}{(v+kt)^{\frac{a}{k}}} dt, \quad |\arg v| < \pi, Re(u) > 0.$$
(2.2)

Proof. Consider

$$D_k \left(\begin{array}{c} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{array} \right) = v^{\frac{-a}{k}} \int_0^\infty t^{u-1} (1-k\frac{t}{v})^{\frac{\delta}{k}-1} e^{-p\frac{t^k}{k}} F_k[(a;k), (b;k); (c;k), -\frac{t}{v}] dt.$$

Now by putting p = 1, $\delta = k$ and b = c, we get

$$D_k \left(\begin{array}{c} (a;k), (b;k); (b;k), 1\\ u, v, k \end{array} \right) = v^{\frac{-a}{k}} \int_0^\infty t^{u-1} e^{-\frac{t^k}{k}} F_k[(a;k), (b;k); (b;k), -\frac{t}{v}] dt.$$
(2.3)

Now since $F_k[(a;k), (b;k); (b;k), -\frac{t}{v}] = (1 + k\frac{t}{v})^{\frac{-a}{k}}$. By substituting this result in (2.3), we have

$$D_{k}\left(\begin{array}{c}(a;k),(b;k);(b;k),1\\u,v,k\end{array}\right) = v^{\frac{-a}{k}} \int_{0}^{\infty} t^{u-1} e^{-\frac{t^{k}}{k}} (1+k\frac{t}{v})^{\frac{-a}{k}} dt$$
$$= \int_{0}^{\infty} \frac{t^{u-1} e^{-\frac{t^{k}}{k}}}{(v+kt)^{\frac{a}{k}}} dt$$
$$= \Gamma_{a,k}(u,v),$$

which is an improved version of Kobayashi's gamma function. The substitution of k = 1 in (2.2) will reduces to Kobayashi's gamma function $\Gamma_a(u, v)$. **Remark.** If a = 0, then (2.2) becomes gamma k-function which is given by

$$\Gamma_k(u) = \int_0^\infty t^{u-1} e^{-\frac{t^k}{k}} dt.$$
(2.4)

Similarly if a = 0 and k = 1, then (2.2) becomes the classical gamma function defined by (1.2).

Proposition 2.1. Let a, b, c and p be complex numbers, then D can be represented as

$$D_k \left(\begin{array}{c} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{array} \right) = v^{u-\frac{a}{k}} \int_0^\infty s^{u-1} (1-ks)^{\frac{\delta}{k}-1} e^{-p\frac{(vs)^k}{k}} F_k[(a,k), (b,k); (c,k); -s] ds, (2.5)$$

where $v \in W_v$ and $u \in W_u$.

Proof. Assume that v > 0 and by changing the variable of integration $s = \frac{t}{v}$ in

(2.1), we have

$$\begin{aligned} D_k \left(\begin{array}{c} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{array} \right) &= v^{-\frac{a}{k}} \int_0^\infty (sv)^{u-1} (1-ks)^{\frac{\delta}{k}-1} e^{-p\frac{(vs)^k}{k}} F_k[(a,k), (b,k); (c,k); -s] v ds \\ &= v^{u-\frac{a}{k}} \int_0^\infty (s)^{u-1} (1-ks)^{\frac{\delta}{k}-1} e^{-p\frac{(vs)^k}{k}} F_k[(a,k), (b,k); (c,k); -s] ds, \end{aligned}$$

which completes the proof.

Theorem 2.2 Let a, b, c and p be complex numbers, then

$$\frac{d^n}{du^n}D_k = v^{\frac{-a}{k}} \int_0^\infty t^{u-1}(1-k\frac{t}{v})^{\frac{\delta}{k}-1}e^{-p\frac{t^k}{k}}(\ln t)^n F_k[(a;k),(b;k);(c;k),-\frac{t}{v}]dt$$

 $\mathbf{Proof} \ \ \mathbf{Consider}$

$$\frac{d^n}{du^n}D_k = v^{\frac{-a}{k}} \int_0^\infty \frac{d^n}{du^n} t^{u-1} (1-k\frac{t}{v})^{\frac{\delta}{k}-1} e^{-p\frac{t^k}{k}} F_k[(a;k),(b;k);(c;k),-\frac{t}{v}] dt.$$
(2.6)

Now since $\frac{d^n}{du^n}t^{u-1} = t^{u-1}(\ln t)^n$, thus by substituting this result in (2.6), we get the required result.

3. The Generalized Incomplete and Complementary Incomplete Functions of D_k .

In this section, we introduce the generalized incomplete and complementary incomplete gamma k-functions of $D_k \begin{pmatrix} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{pmatrix}$. If $\omega > 0$, then we defined the generalized incomplete gamma k-function as

$$D_{0,k}^{\omega} \begin{pmatrix} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{pmatrix} = v^{\frac{-a}{k}} \int_{0}^{\omega} t^{u-1} (1-k\frac{t}{v})^{\frac{\delta}{k}-1} e^{-p\frac{t^{k}}{k}} F_{k}[(a;k), (b;k); (c;k), -\frac{t}{v}] dt, \quad (3.1)$$

where $v \in W_v = \{v \in \mathbb{C} : |\arg v| < \pi\}, u \in W_u = \{u \in \mathbb{C} : Re(u) > 0\}, c$ are complex parameters with $c \neq 0, 1, 2, \cdot$ and $F_k[(a; k), (b; k); (c; k), -\frac{t}{v}]$ is the Gauss hypergeometric k-functions, and the generalized complementary incomplete gamma k-function is given by

$$D_{\omega,k}^{\infty} \begin{pmatrix} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{pmatrix} = v^{\frac{-a}{k}} \int_{\omega}^{\infty} t^{u-1} (1-k\frac{t}{v})^{\frac{\delta}{k}-1} e^{-p\frac{t^{k}}{k}} F_{k}[(a;k), (b;k); (c;k), -\frac{t}{v}] dt, \quad (3.2)$$

with the same conditions given in (3.1). Thus, this implies that

$$D_k \left(\begin{array}{c} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{array} \right) = D_{0,k}^{\omega} \left(\begin{array}{c} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{array} \right) + D_{\omega,k}^{\infty} \left(\begin{array}{c} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{array} \right).$$

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For $\delta = k, p = 1, b = c$ and a = m where $m \in \mathbb{Z}^+$ and using the relation for Gauss hypergeometric function [6]

$$F_k[(a;k),(b;k);(b;k),z] = (1-kz)^{\frac{-a}{k}}, \quad \arg|1-kz| < \pi,$$

then (3.2) reduces to the generalized incomplete gamma k-function discussed in section 2 .

$$\Gamma_{m,k}(u,v,\omega) = \int_{\omega}^{\infty} \frac{t^{u-1}e^{-\frac{t^k}{k}}}{(v+kt)^{\frac{m}{k}}} dt.$$
(3.3)

If m = 0, then this result reduces to the well known complementary incomplete gamma k-function [6]

$$\Gamma_k(u,\omega) = \int_{\omega}^{\infty} t^{u-1} e^{-\frac{t^k}{k}} dt,$$
(3.4)

and if k = 1 this result reduces to the classical incomplete gamma function

$$\Gamma(u,\omega) = \int_{\omega}^{\infty} t^{u-1} e^{-t} dt.$$

On the other hand if $\delta = k$, p = 1, b = c and a = 0, then (3.1) reduces to the classical incomplete gamma function

$$\gamma(u,\omega) = \int_{0}^{\omega} t^{u-1} e^{-t} dt.$$

4. Relations Involving The Function D_k

In this section, we derive some recurrence relations which containing more generalized gamma k-function $D_k \begin{pmatrix} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{pmatrix}$. The function $D_k \begin{pmatrix} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{pmatrix}$ is not symmetric with respect to the parameters a and b. In fact,

$$D_k \left(\begin{array}{c} (b;k), (a;k); (c;k), p \\ u, v, \delta \end{array} \right) = v^{\frac{a-b}{k}} D_k \left(\begin{array}{c} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{array} \right).$$
(4.1)

Proposition 3.1. The following relations hold for k > 0;

$$(b-a)D_k \left(\begin{array}{c} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{array} \right) + avD_k \left(\begin{array}{c} (a+k;k), (b;k); (c;k), p \\ u, v, \delta \end{array} \right) \\ - bD_k \left(\begin{array}{c} (a;k), (b+k;k); (c;k), p \\ u, v, \delta \end{array} \right) = 0, \quad (4.2)$$

$$(c-a-k)D_{k}\left(\begin{array}{c}(a;k),(b;k);(c;k),p\\u,v,\delta\end{array}\right) + avD_{k}\left(\begin{array}{c}(a+k;k),(b;k);(c;k),p\\u,v,\delta\end{array}\right) \\ - (c-k)D_{k}\left(\begin{array}{c}(a;k),(b;k);(c-k;k),p\\u,v,\delta\end{array}\right) = 0, \quad (4.3)$$

$$(c-a-b)D_k \left(\begin{array}{c} (a;k), (b;k); (c;k), p \\ u,v,\delta \end{array} \right) - (c-b)D_k \left(\begin{array}{c} (a;k), (b-k;k); (c;k), p \\ u,v,\delta \end{array} \right) + avD_k \left(\begin{array}{c} (a+k;k), (b;k); (c;k), p \\ u,v,\delta \end{array} \right) + akD_k \left(\begin{array}{c} (a+k;k), (b;k); (c;k), p \\ u+1,v,\delta \end{array} \right) = 0,$$

$$(4.4)$$

$$(c-b-a)D_k \left(\begin{array}{c} (a;k), (b;k); (c;k), p\\ u,v \end{array}\right) - (c-a)D_k \left(\begin{array}{c} (a-k;k), (b;k); (c;k), p\\ u,v,\delta \end{array}\right) + bvD_k \left(\begin{array}{c} (a;k), (b+k;k); (c;k), p\\ u,v,\delta \end{array}\right) + bkD_k \left(\begin{array}{c} (a;k), (b+k;k); (c;k), p\\ u+1,v,\delta \end{array}\right) = 0,$$

$$(4.5)$$

$$(b-a)vD_k \left(\begin{array}{c} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{array} \right) - (c-a)D_k \left(\begin{array}{c} (a-k;k), (b;k); (c;k), p \\ u, v, \delta \end{array} \right) \\ + (b-a)kD_k \left(\begin{array}{c} (a;k), (b;k); (c;k), p \\ u+1, v, \delta \end{array} \right) + (c-b)vD_k \left(\begin{array}{c} (a;k), (b-k;k); (c;k), p \\ u, v\delta \end{array} \right) = 0,$$

$$(4.6)$$

$$(c-b-k)D_{k}\left(\begin{array}{c}(a;k),(b;k);(c;k),p\\u,v,\delta\end{array}\right) + bvD_{k}\left(\begin{array}{c}(a;k),(b+k;k);(c;k),p\\u,v,\delta\end{array}\right) - (c-k)D_{k}\left(\begin{array}{c}(a;k),(b;k);(c-k;k),p\\u,v,\delta\end{array}\right) = 0.$$
(4.7)

Proof. To prove the relation (4.2), we use (1.9) by taking $F_k[a, b; c, z] = F_k[(a; k), (b; k); (c; k), z]$ as

 $(b-a)F_k[(a;k), (b;k); (c;k), z] + aF_k[(a+k;k), (b;k); (c;k), z] - bF_k[(a;k), (b+k;k); (c;k), z] = 0.$ By putting $z = \frac{-t}{v}$ in above relation, we have

$$\begin{split} (b-a)D_k\left(\begin{array}{c}(a;k),(b;k);(c;k),p\\u,v,\delta\end{array}\right)\\ &= v^{\frac{-a}{k}}\int_0^\infty t^{u-1}(1-k\frac{t}{v})^{\frac{\delta}{k}-1}e^{\frac{-pt^k}{k}}\times [-aF_k[(a+k,k),(b,k);(c,k);-\frac{t}{v}]\\ &+ bF_k[(a,k),(b+k,k);(c,k);-\frac{t}{v}]]dt\\ &= -a[v^{-(\frac{a}{k}-1)}\int_0^\infty t^{u-1}(1-k\frac{t}{v})^{\frac{\delta}{k}-1}e^{\frac{-pt^k}{k}}\times F_k[(a+k;k),(b;k);(c;k),-\frac{t}{v}]]dt\\ &+ bv^{\frac{-a}{k}}\int_0^\infty t^{u-1}(1-k\frac{t}{v})^{\frac{\delta}{k}-1}e^{\frac{-pt^k}{k}}F_k[(a;k),(b+k;k);(c;k),-\frac{t}{v}]dt\\ &= -avD_k\left(\begin{array}{c}(a+k;k),(b;k);(c;k),p\\u,v\end{array}\right)+bD_k\left(\begin{array}{c}(a;k),(b+k;k);(c;k),p\\u,v\end{array}\right),\end{split}$$

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which is the required relation. On the similar way we can prove relation (4.3). To prove relation (4.4), we use (1.11) as;

$$(a+b-c)F_k[(a;k),(b;k);(c;k);z] = a(1-kz)F_k[(a+k;k),(b;k);(c;k);z) - (c-b)F_k[(a;k),(b-k;k);(c;k);z].$$

with $z = -\frac{t}{v}$. Now, we have

$$(c-a-b)D_k\left(\begin{array}{c}(a;k),(b;k);(c;k),p\\u,v,\delta\end{array}\right)$$

$$\begin{split} &= v^{\frac{-a}{k}} \int_{0}^{\infty} t^{u-1} (1-k\frac{t}{v})^{\frac{\delta}{k}-1} e^{\frac{-pt^{k}}{k}} \times (c-a-b) F_{k}[(a+k,k),(b,k);(c,k);-\frac{t}{v}] dt \\ &= v^{\frac{-a}{k}} \int_{0}^{\infty} t^{u-1} (1-k\frac{t}{v})^{\frac{\delta}{k}-1} e^{\frac{-pt^{k}}{k}} (c-b) F_{k}[(a,k),(b-k,k);(c,k);-\frac{t}{v}] \\ &- a(1+k\frac{t}{v}) bF_{k}[(a+k,k),(b,k);(c,k);-\frac{t}{v}]] dt \\ &= (c-b)[v^{\frac{-a}{k}} \int_{0}^{\infty} t^{u-1} (1-k\frac{t}{v})^{\frac{\delta}{k}-1} e^{\frac{-pt^{k}}{k}} \times F_{k}[(a;k),(b-k;k);(c;k),-\frac{t}{v}]] dt \\ &+ av^{-(\frac{a}{k}-1)} \int_{0}^{\infty} t^{u-1} (1-k\frac{t}{v})^{\frac{\delta}{k}-1} e^{\frac{-pt^{k}}{k}} F_{k}[(a+k;k),(b;k);(c;k),-\frac{t}{v}] dt \\ &- akv^{-(\frac{a}{k})} \int_{0}^{\infty} t^{u} (1-k\frac{t}{v})^{\frac{\delta}{k}-1} e^{\frac{-pt^{k}}{k}} F_{k}[(a+k;k),(b;k);(c;k),-\frac{t}{v}] dt \\ &= (c-b)vD_{k} \left(\begin{array}{c} (a;k),(b-;k);(c;k),p \\ u,v,\delta \end{array} \right) - avD_{k} \left(\begin{array}{c} (a+k;k),(b;k);(c;k),p \\ u,v,\delta \end{array} \right) \\ &- akD_{k} \left(\begin{array}{c} (a+k;k),(b;k);(c;k),p \\ u+1,v,\delta \end{array} \right) , \end{split}$$

which completes the proof. On the similar method, we can prove relation (4.5), (4.6) and (4.7) by using (1.12), (1.13) and (1.14) respectively.

5. Asymptotic Expansions for Generalized Gamma k-Functions

In this section, we introduce the asymptotic expansion for $D_k \begin{pmatrix} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{pmatrix}$, $D_{0,k}^{\omega} \begin{pmatrix} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{pmatrix}$ and $D_{\omega,k}^{\infty} \begin{pmatrix} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{pmatrix}$. We assume that $v \to \infty$ and using the expansion in series of the Gauss hypergeometric k-function

$${}_{2}F_{1,k}[a,b;c;z] = \sum_{n=0}^{\infty} \frac{(a)_{n,k}(b)_{n,k}}{(c)_{n,k}} \frac{z^{n}}{n!}, \quad |z| < 1, k > 0,$$

where $z = -\frac{t}{v}$. After substitution in (2.1) and interchanging the order of integral and sum, we obtain

$$D_k \left(\begin{array}{c} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{array} \right) = v^{-\frac{a}{k}} \sum_{n=0}^{\infty} \frac{(-1)^n (a)_{n,k} (b)_{n,k}}{(c)_{n,k} n! v^n} \int_0^{\infty} t^{n+u-1} (1-k\frac{t}{v})^{\frac{\delta}{k}-1} e^{-p\frac{t^k}{k}} dt.$$
(5.1)

Now since

$$(1-kz)^{-\frac{a}{k}} = \sum_{m=0}^{\infty} \frac{(a)_{m,k}}{m!} z^m, \quad |z| < 1$$

by substituting in (5.1), we have

$$D_k \left(\begin{array}{c} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{array} \right) = v^{-\frac{a}{k}} \sum_{n=0}^{\infty} \frac{(-1)^n (a)_{n,k} (b)_{n,k}}{(c)_{n,k} n! v^n} \sum_{m=0}^{\infty} \frac{(k-\delta)_{m,k}}{m! v^m} \\ \times \int_0^{\infty} t^{n+m+u-1} e^{-p \frac{t^k}{k}} dt.$$

Substituting $y = pt^k$ and $dt = \frac{1}{kp^{\frac{1}{k}}}y^{\frac{1}{k}-1}$ in above equation, we get

$$D_k \left(\begin{array}{c} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{array} \right) = \frac{v^{-\frac{a}{k}}}{p^{\frac{u}{k}}} \sum_{n=0}^{\infty} \frac{(-1)^n (a)_{n,k} (b)_{n,k}}{(c)_{n,k} n!} \sum_{m=0}^{\infty} \frac{(k-\delta)_{m,k}}{m! p^{\frac{m+n}{k}} v^{n+m}} \\ \times \frac{1}{k} \int_0^\infty y^{\frac{n+m+u}{k} - 1} e^{-\frac{y}{k}} dy. \quad (5.2)$$

Again by variable substitution the integral part becomes

$$\frac{1}{k} \int_{0}^{\infty} y^{\frac{n+m+u}{k}-1} e^{-\frac{y}{k}} dy = \int_{0}^{\infty} t^{n+m+u-1} e^{-\frac{t^{k}}{k}} dt$$
$$= \Gamma_{k}(n+m+u).$$

Using this result in (5.2), we get

$$D_k \left(\begin{array}{c} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{array} \right) = \frac{v^{-\frac{a}{k}}}{p^{\frac{u}{k}}} \sum_{n=0}^{\infty} \frac{(-1)^n (a)_{n,k} (b)_{n,k}}{(c)_{n,k} n! v^n p^{\frac{n}{k}}} \sum_{m=0}^{\infty} \frac{(k-\delta)_{m,k}}{m! p^{\frac{m}{k}} v^m} \Gamma_k (n+m+u),$$
(5.3)

which is the required asymptotic expansion for the generalized gamma k-function. Following the same procedure, we obtain the asymptotic expansion of generalized incomplete gamma k-function from definition (3.1)

$$D_{k}\left(\begin{array}{c}(a;k),(b;k);(c;k),p\\u,v,\delta\end{array}\right) = \frac{v^{-\frac{a}{k}}}{p^{\frac{u}{k}}}\sum_{n=0}^{\infty}\frac{(-1)^{n}(a)_{n,k}(b)_{n,k}}{(c)_{n,k}n!}\sum_{m=0}^{\infty}\frac{(k-\delta)_{m,k}}{m!p^{\frac{m+n}{k}}v^{n+m}} \\ \times \int_{0}^{p\omega}t^{n+m+u-1}e^{-\frac{t^{k}}{k}}dt, \quad (5.4)$$

and by using the result (5.5), this can be written as

$$D_{0,k}^{\omega} \left(\begin{array}{c} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{array} \right) = \frac{v^{-\frac{a}{k}}}{p^{\frac{u}{k}}} \sum_{n=0}^{\infty} \frac{(-1)^n (a)_{n,k} (b)_{n,k}}{(c)_{n,k} n! v^n p^{\frac{n}{k}}} \sum_{m=0}^{\infty} \frac{(k-\delta)_{m,k}}{m! p^{\frac{m}{k}} v^m} \gamma_k (n+m+u, p\omega).$$
(5.5)

Similarly, we can find the asymptotic expansion of complementary incomplete gamma k-function. From definition (3.2), we have

$$D_{\omega,k}^{\infty} \left(\begin{array}{c} (a;k), (b;k); (c;k), p \\ u, v, \delta \end{array} \right) = \frac{v^{-\frac{a}{k}}}{p^{\frac{u}{k}}} \sum_{n=0}^{\infty} \frac{(-1)^n (a)_{n,k} (b)_{n,k}}{(c)_{n,k} n!} \sum_{m=0}^{\infty} \frac{(k-\delta)_{m,k}}{m! p^{\frac{m+n}{k}} v^{n+m}} \\ \times \int_{p\omega}^{\infty} t^{n+m+u-1} e^{-\frac{t^k}{k}} dt$$

$$= \frac{v^{-\frac{a}{k}}}{p^{\frac{u}{k}}} \sum_{n=0}^{\infty} \frac{(-1)^n (a)_{n,k} (b)_{n,k}}{(c)_{n,k} n! v^n p^{\frac{n}{k}}} \sum_{m=0}^{\infty} \frac{(k-\delta)_{m,k}}{m! p^{\frac{m}{k}} v^m} \Gamma_k(n + m + u, p\omega).$$
(5.6)

6. CONCLUSION AND SUMMARY

In this paper, the aim of the authors is to consider further generalization of the gamma k-function. Also, to obtain some properties, recurrence relations involving generalized gamma k-function and establish the asymptotic expansions for the said functions in term of k where k > 0. Obviously the substitution of k = 1 will lead to the classical results of the said functions.

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